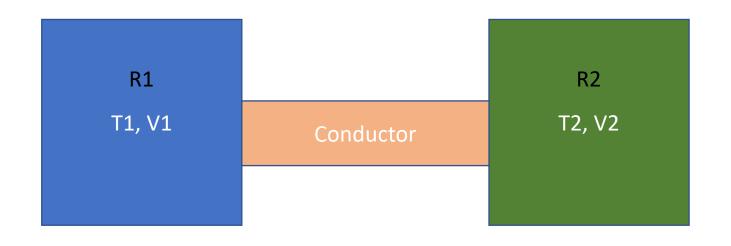
Tutorial #6:
Irreversible thermodynamics



Because of the gradient in temperature and voltage, we can express the flux between the two reservoirs:

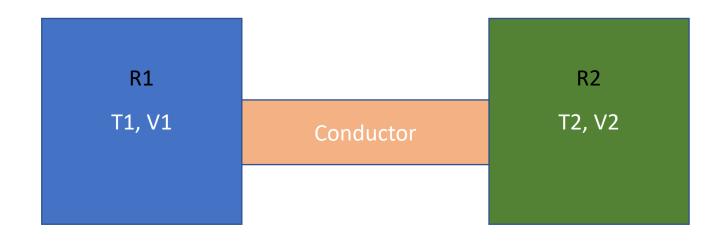
$$I_q = l_{11}\Delta V + l_{12}\Delta T$$

$$I_U = l_{21}\Delta V + l_{22}\Delta T$$

But with this formulation, there is no reciprocity of the coefficients: I21 = L12 is NOT satisfied

In this exercise, we will define adequate quantities and derive coefficients satisfying recirpocity relations.

Generalized force + exchange of an extensive physical quantity.



1) Write an energy and charge budget between the two réservoirs:

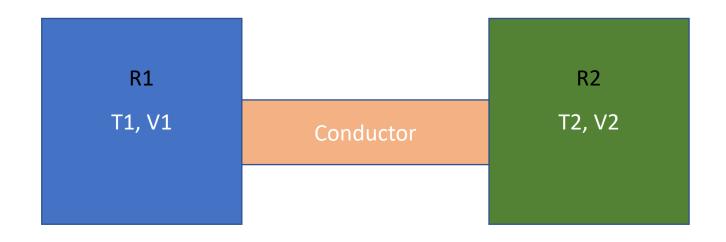
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$$I_q = l_{11}\Delta V + l_{12}\Delta T$$

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2) Derive the variation of entropy in R1 and R2. You can use the thermodynamic identity:

$$dU = TdS + \mu dN + Vdq$$



1) Write an energy and charge budget between the two réservoirs:

$$q\delta N_1 = -q\delta N_2$$
$$\delta U_1 = -\delta U_2$$

Because of the gradient in temperature and voltage, we can express the flux between the two reservoirs:

$$I_q = l_{11}\Delta V + l_{12}\Delta T$$

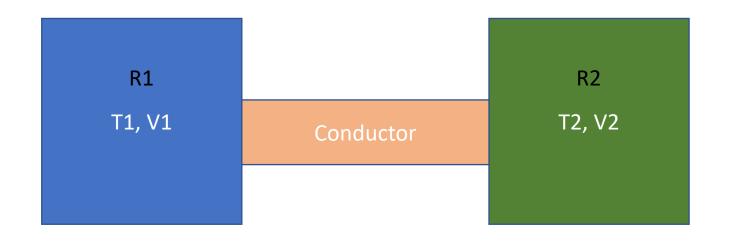
$$I_U = l_{21}\Delta V + l_{22}\Delta T$$

2) Derive the variation of entropy in R1 and R2. You can use the thermodynamic identity:

$$\delta S_{1} = \frac{\delta U_{1}}{T_{1}} - \frac{\mu(T_{1}) - qV_{1}}{T_{1}} \delta N_{1}$$

$$= -\frac{\delta U_{2}}{T_{1}} + \frac{\mu(T_{1}) - qV_{1}}{T_{1}} \delta N_{2}$$

$$\delta S_{2} = \frac{\delta U_{2}}{T_{2}} - \frac{\mu(T_{2}) + qV_{2}}{T_{2}} \delta N_{2}$$

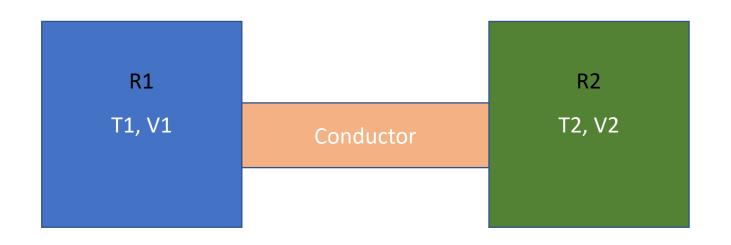


Because of the gradient in temperature and voltage, we can express the flux between the two reservoirs:

$$I_q = l_{11}\Delta V + l_{12}\Delta T$$

$$I_U = l_{21}\Delta V + l_{22}\Delta T$$

3) Derive the total variation of entropy (i.e. variation in both réservoirs)



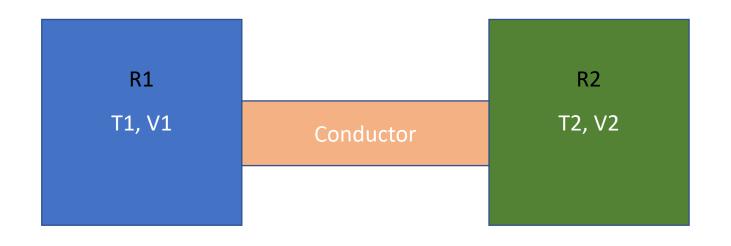
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3) Derive the total variation of entropy (i.e. variation in both réservoirs)

$$\begin{split} \delta S &= \delta S_1 + \delta S_2 \\ &= (\frac{1}{T_2} - \frac{1}{T_1}) \delta U_2 + (\frac{\mu(T_1) + qV_1}{T_1} - \frac{\mu(T_2) + qV_2}{T_2}) \delta N_2 \end{split}$$



Because of the gradient in temperature and voltage, we can express the flux between the two reservoirs:

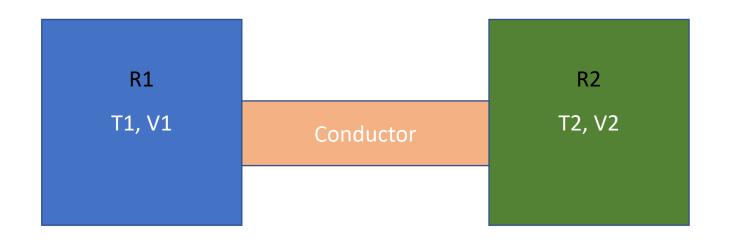
$$I_q = l_{11}\Delta V + l_{12}\Delta T$$

$$I_U = l_{21}\Delta V + l_{22}\Delta T$$

4) Derive the generalized forces Xu and Xq associated to the exchange of energy and charge, knowing that the rate of creation of entropy can be expressed using:

$$\frac{\Delta S}{\Delta t} = \sum_{i} J_i X_i, \quad \text{With rate of extensive variables given by :}$$

$$J_U = \frac{\Delta U}{\Delta t} \quad J_q = q \frac{\Delta N}{\Delta t}.$$



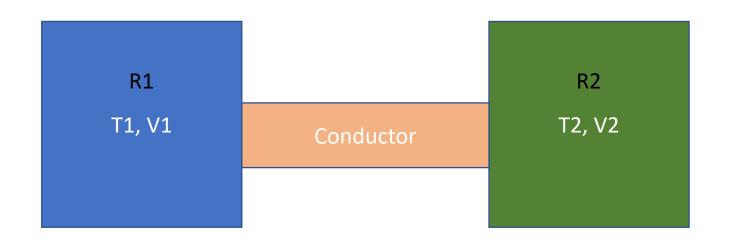
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4) Derive the generalized forces Xu and Xq associated to the exchange of energy and charge:

$$\frac{\delta S}{\delta t} = \left(\frac{1}{T_2} - \frac{1}{T_1}\right) \underbrace{\frac{\delta U_2}{\delta t}}_{J_U} + \left(\frac{\mu(T_1)/q + V_1}{T_1} - \frac{\mu(T_2)/q + V_2}{T_2}\right) \underbrace{q\frac{\delta N_2}{\delta t}}_{J_q}$$



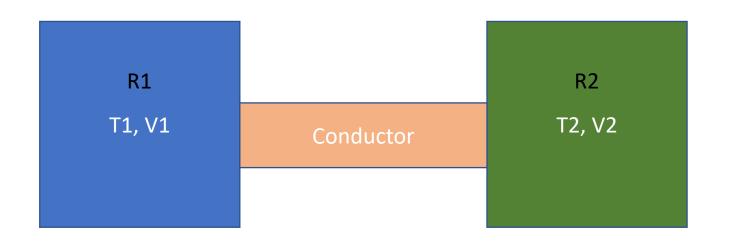
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4) Derive the generalized forces Xu and Xq associated to the exchange of energy and charge:

$$\begin{split} \frac{\delta S}{\delta t} &= (\frac{1}{T_2} - \frac{1}{T_1}) \underbrace{\frac{\delta U_2}{\delta t}}_{J_{TI}} + (\frac{\mu(T_1)/q + V_1}{T_1} - \frac{\mu(T_2)/q + V_2}{T_2}) \underbrace{q \underbrace{\frac{\delta N_2}{\delta t}}_{J_q}}_{J_q} \\ &= \Delta (\frac{1}{T}) J_U - \frac{1}{q} \Delta (\frac{\mu + qV}{T}) J_q \end{split}$$



Because of the gradient in temperature and voltage, we can express the flux between the two reservoirs:

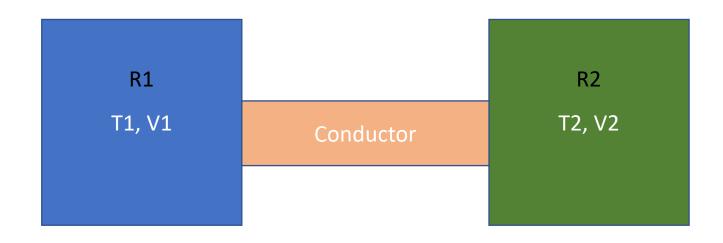
$$I_q = l_{11}\Delta V + l_{12}\Delta T$$

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4) Derive the generalized forces Xu and Xq associated to the exchange of energy and charge:

$$\frac{\delta S}{\delta t} = \Delta(\frac{1}{T})J_U - \frac{1}{q}\Delta(\frac{\mu + qV}{T})J_q$$

$$X_U = \Delta(\frac{1}{T}) = -\frac{\Delta T}{T^2} \qquad X_q = -\frac{1}{q}\Delta(\frac{\mu_{ec}}{T})$$



We know now how to define coefficients that satisfy Onsager reciprocity

5) Derive the thermal conductance given by $\ J_{II} = -\Gamma \Delta T$

Because of the gradient in temperature and voltage, we can express the flux between the two reservoirs:

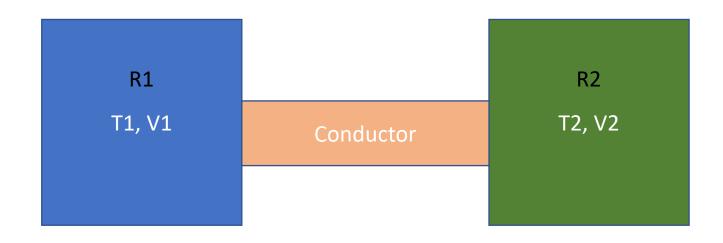
$$I_{q} = l_{11}\Delta V + l_{12}\Delta T$$

$$I_{U} = l_{21}\Delta V + l_{22}\Delta T$$

$$J_{q} = L_{11}X_{q} + L_{12}X_{U}$$

$$J_{U} = L_{21}X_{q} + L_{22}X_{U}$$

Which is given in absence of charge flux (instead of zero voltage)!



We know now how to define coefficients that satisfy Onsager reciprocity

$$J_U = -\Gamma \Delta T$$
$$= L_{21}X_q + L_{22}X_U$$

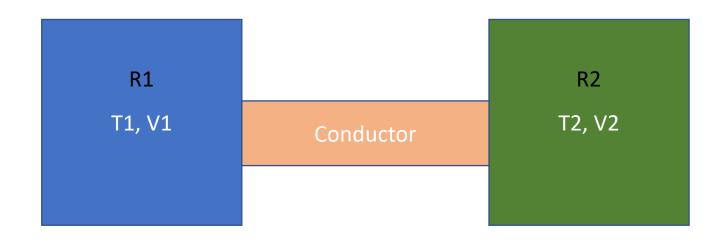
Because of the gradient in temperature and voltage, we can express the flux between the two reservoirs:

$$I_{q} = l_{11}\Delta V + l_{12}\Delta T$$

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$$J_{q} = L_{11}X_{q} + L_{12}X_{U}$$

$$J_{U} = L_{21}X_{q} + L_{22}X_{U}$$



We know now how to define coefficients that satisfy Onsager reciprocity

$$J_U = -\Gamma \Delta T$$
$$= L_{21}X_q + L_{22}X_U$$

$$J_q=0$$

$$X_q=-\frac{L_{12}}{L_{11}}X_U \quad \text{with} \quad X_U=\Delta(\frac{1}{T})=-\frac{\Delta T}{T^2}$$

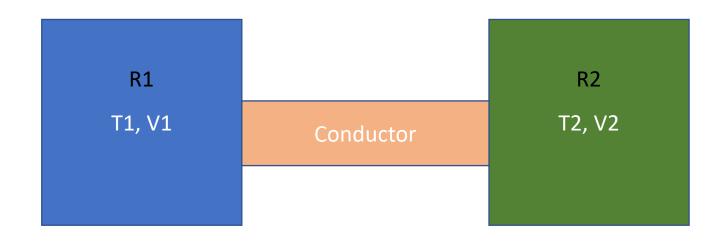
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We know now how to define coefficients that satisfy Onsager reciprocity

$$J_U = -\Gamma \Delta T$$

$$= L_{21} X_q + L_{22} X_U$$

$$J_q = 0$$

$$X_q = -\frac{L_{12}}{L_{11}} X_U \quad \text{with} \quad X_U = \Delta(\frac{1}{T}) = -\frac{\Delta T}{T^2}$$

Because of the gradient in temperature and voltage, we can express the flux between the two reservoirs:

$$I_{q} = l_{11}\Delta V + l_{12}\Delta T$$

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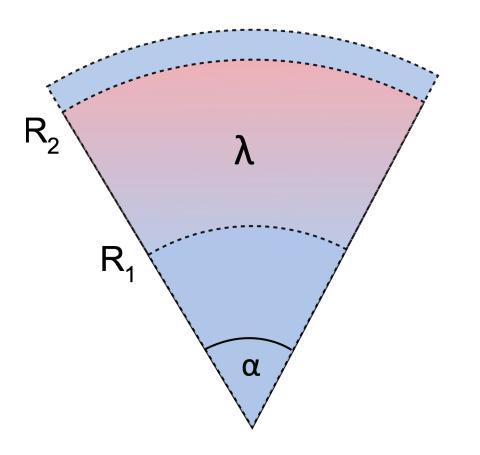
$$J_{U} = L_{21}X_{q} + L_{22}X_{U}$$

$$-\Gamma \Delta T = -\frac{L_{21}L_{12}}{L_{11}}X_U + L_{22}X_U$$

$$\Gamma = \frac{L_{21}L_{12} - L_{22}L_{11}}{L_{11}} \frac{-1}{T^2}$$

$$= \frac{L_{22}L_{11} - L_{12}^2}{L_{11}T^2}$$

Exercise #2: Heat exchange in a constriction



$$\mathbf{J}_{\mathbf{U}} = -\lambda \vec{\nabla} T = -\lambda \frac{\partial T}{\partial r} \mathbf{u_r}$$

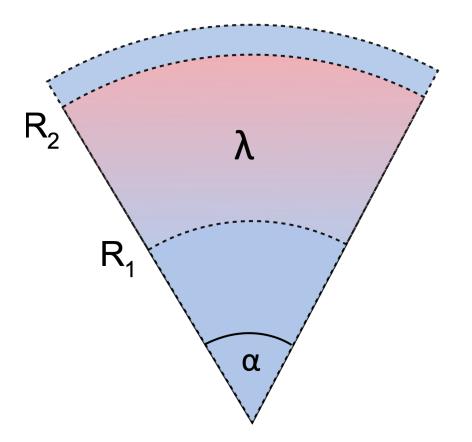
 $\Omega = 2\pi(1 - \cos\alpha)$

$$\Sigma(r) = \Omega r^2 = 2\pi (1 - \cos \alpha) r^2$$

Radial flux of heat across « shell portions » defined by Sigma(r)

- Rate of local volumic entropy creation ? (remember it is affinity x flux)
- 2) Show that it is proportionnal to 1/r⁴ (use the fact that the flux is constant across Sigma(r))
- 3) Where do irreversible phenomena occur most in the system?

Exercise #2: Heat exchange in a constriction



$$\mathbf{J_U} = -\lambda \vec{\nabla} T = -\lambda \frac{\partial T}{\partial r} \mathbf{u_r}$$

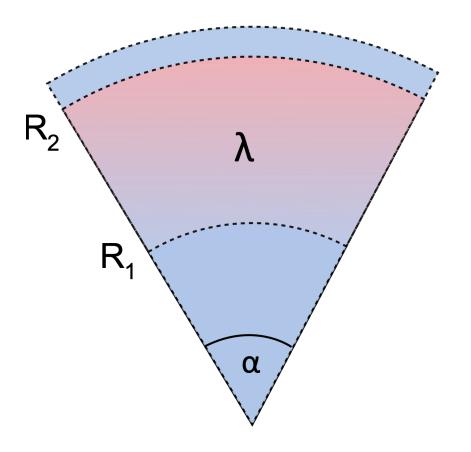
$$\Omega = 2\pi (1 - \cos \alpha)$$

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1) Rate of local volumic entropy creation ? (remember it is affinity x flux)

$$\frac{\partial s}{\partial t} = \vec{\nabla} (\frac{1}{T}) \cdot \mathbf{J}_{\mathbf{U}}$$
$$= \frac{\lambda}{T^2} \vec{\nabla} T \cdot \vec{\nabla} T$$
$$= \frac{\lambda}{T^2} \left(\frac{\partial T}{\partial r} \right)^2$$

Exercise #2: Heat exchange in a constriction



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$$\Omega = 2\pi (1 - \cos \alpha)$$

$$\Sigma(r) = \Omega r^2 = 2\pi (1 - \cos \alpha) r^2$$

2) Show that it is proportionnal to $1/r^4$ (use the fact that the flux is constant across Sigma(r))

$$\mathbf{J_u}(r) \cdot \mathbf{\Sigma}(r) = \text{cste}$$

$$-\lambda \frac{\partial T}{\partial r} 2\pi (1 - \cos \alpha) r^2 = \text{cste}$$

$$\frac{\partial T}{\partial r} = \frac{A}{r^2}$$

$$\frac{\partial s}{\partial t} = \frac{\lambda}{T^2} \left(\frac{A}{r^2}\right)^2 \approx \frac{1}{r^4}$$